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Plasma Densities and Irregularities at 830 km Altitude Based on Observations During 1979

FREDERICK J. RICH MICHAEL SMIDDY



28 May 1986



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SPACE PHYSICS DIVISION

PROJECT 2311

AIR FORCE GEOPHYSICS LABORATORY

HANSCOM AFB, MA 01731

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Plasma Densities and Irregularities at 830 km Altitude Based on Observations During 1979

1. INTRODUCTION

Ground-based profiles of the ionospheric density have been obtained in vast quantities for over half a century. However, information about the ionospheric profile above the maximum density (f_0F_2) of the ionosphere (at 325 to 475 km altitude) can only be obtained from satellites and incoherent scatter radar measurements which are much more limited in quantity. As a result, the standard models of the ionosphere, such as IRI 79 (Rawer et al and Rawer) have made extensive use of the ground-based data and have made much less use of satellite data. The limited topside ionosphere data which has been used has been weighted more for the f_0F^2 altitude than for higher altitudes. We have found these models to be inadequate when compared to topside density measurements at 830 km altitudes.

Thermal plasma density measurements in the topside ionosphere have been made with many satellites including ESRO 1A, ESRO 4, OGO 6, Alouette, ISIS I, ISIS II, Atmospheric Explorer C and D and ISS-b. While much data from these satellites have been published for particular seasons or latitude or longitudes, little

⁽Received for publication 28 May 1986)

Rawer, K., Bilitza, D., and Ramakrishnan, S. (1978) Goals and status of the international reference ionosphere, <u>Rev. Geophys. Sp. Phys.</u>, 16(No. 2):177-181.

Rawer, K. (1981) International reference ionosphere - IRI 79, Report UAG-82, World Data Center - A, NOAA, Boulder, Colorado.

of these data have been compiled for wide spans of latitude, longitude and seasons. The best compilation to date is the ISS-b data set (Hakura, and Maruyama and Matuura). To add to the database from which a more realistic model of both the topside and bottomside ionosphere can be constructed, we have compiled all of the density data from the F2 and F4 satellites during 1979 for all latitudes (except polar regions), all longitudes and all seasons.

Since regions of irregularity in the plasma density are often associated with anomalous plasma densities, many observations of the plasma density have been excluded from the survey. However, since regions of disturbed plasma are often related to radio scintillation we also report on the probability of finding disturbed plasma, especially in the equatorial region.

The disturbed plasma near the equator is often associated with rising plumes of low density plasma in the evening sector breaking through the bottomside of the F layer and floating to the topside ionosphere. This phenomenon is generally referred to as spread-F due to its signature on ionograms. The frequency of occurrence of disturbed plasma as a function of season, latitude and longitude has been noted and is compared with previous observations. While the plasma scale lengths observed by the DMSP instrument (1 to 35 km) is greater than the scale required for radio disturbance, it is reasonable to assume that the plasma within the plumes is turbulent over a wide range of scale length up to the width of the plume.

The disturbed plasma at high latitudes is associated with the nighttime trough, the regions of precipitating particles, especially 0.1 to 30 keV electrons, and to clouds or "blobs" of plasma density enhancement which are convected over the pole from the dayside to the nightside. These processes tend to create steep plasma density gradients with scale lengths of several km. Other process act on these steep gradients to cause plasma density of small enough scale lengths to cause degradation of radio signals. Basu et al⁵ reported upon a simultaneous DMSP observation of precipitating electrons and a steep plasma density gradient and a ground observation of radio scintillation.

^{3.} Hakura, Y. (Ed.) Atlas of proton, helium ion and oxygen ion densities obtained from Ionospheric Sounding Satellite-b in situ observations. October 1978 to August 1979 Radio Research Lab., Ministry of Posts and Telecommunications, Tokyo, Japan.

Maruyama, T., and Matuura, N. (1984) Longitudinal variability of annual changes in activity of equatorial spread F and plasma bubbles, J. Geophys. Res., 89(No. A12):10,903-10,912.

Basu, A., MacKenzie, E., Basu, S., Carlson, H.C., Hardy, D.A., Rich, F.J., and Livingston, R.C. (1983) Coordinated measurements of low-energy electron precipitation and scintillations/TEC in the auroral oval, Radio Sci., 18:1151-1165.

2. INSTRUMENTATION

The topside ionospheric plasma monitor (SSIE) has been flown on several satellites of the Defense Meteorological Satellite Program (DMSP) which are in polar, circular, sun-synchronous orbits. The SSIE consists of a gridded, spherical Langmuir probe for measuring thermal electrons and a planar Faraday cup configured as a retarding potential analyzer (RPA) for measuring thermal ions. The diameter of the outer grid of the Langmuir probe is 2.50 in. (5.72 cm). The aperture diameter of the planar RPA is 1.00 in. (2.54 cm). Both instruments are mounted on a 27 in. (69 cm) boom. These instruments have been described by Smiddy et al⁶ and Rich et al⁷ and results have been published by Burke et al⁸ and Young et al. 9 These instruments made measurements of the in situ total ion and electron density seven times per second in the range of 10 to 10⁶ cm⁻³. Any density greater than the maximum of the instrument is reported in the telemetry as the highest possible value $[\log_{10}(\text{density}) = 6.1]$. They also made measurements of the ion and electron temperature once each 64 sec and in the range of 800° to 10,000°K. These instruments use a single range, logarithmic amplifier and a nine-bit analogto-digital conversion. The least significant bit of the telemetry represents a change of density of log, (density) of 0.01. Relative changes (changes from one moment to the next due to local variations) in density of 2 percent are clearly described. The absolute density is a function of the calibration and the spacecraft-plasma interaction. It is estimated to be accurate to better than 25 percent in density or better than ± 0.10 in \log_{10} (density).

In flight, only the total ion density (which is equivalent to the ambient electron density) can be measured whenever the satellite is in sunlight which is approximately 80 percent of the time. This is due to a spacecraft potential caused by unshielded contacts on the solar cell array, which tends to attract electrons and repel ions.

Smiddy, M., Sagalyn, R.C., Sullivan, W.P., Wildman, P.J.L., Anderson, P., and Rich, F. (1978) The Topside Ionosphere Plasma Monitor (SSIE) for the Block 5D Flight 2 DMSP Satellite, Report AFGL-TR-78-0071, AD A058503, Hanscom AFB, Massachusetts.

^{7.} Rich, F., Smiddy, M., Sagalyn, R.C., Burke, W.J., Anderson, P., Bredesen, S., and Sullivan, W.P. (1980) In-flight Characteristics of the Topside Ionospheric Monitor (SSIE) on the DMSP Satellites Flight 2 and Flight 4, Report AFGL-TR-80-0152, AD A088879, Hanscom AFB, Massachusetts.

^{8.} Burke, W.J., Donatelli, D.E., and Sagalyn, R.C. (1980) The longitudinal distribution of equatorial spread F plasma bubbles in the topside ionosphere, J. Geophys. Res., 85:1335-1340.

Young, E. R., Burke, W.a., Rich, F. J., and Sagalyn, R. C. (1984) The distribution of topside spread-F from in situ measurements by Defense Meteorological Satellite Program F2 and F4, J. Geophys. Res., 89(No. A7):5565-5574.

Thus the vehicle is driven to negative potentials of -20 V to -32 V. The measured ion current is unaffected by the positive potential because the spacecraft velocity is much greater than the ion thermal motion, and because ions which are deflected toward the sensor are collected by the external case of the sensor and a ground plane around the ion sensor.

Flight 2 (F2) of the most recent generation of DMSP spacecraft was launched in June 1977 into a 820 km altitude, 99.1° inclination, circular orbit and carried the first of the SSIE instruments. Both the F2 satellite and the SSIE instrument operated continuously until January 1980. The F2 satellite was not quite in a sun-synchronous orbit but in an orbit that shifted approximately one-half hour in local time per year. In 1979, F2 was approximately located with the ascending node at 0930 hours and descending node at 2130 hours geographic local time. Flight 4 (F4) was launched in June 1979 into a 830 km altitude, 98.7° inclination, circular orbit with the second of the SSIE instruments. Both the F4 satellite and the SSIE instrument operated continuously until February 1980. From February 1980 to August 1980, intermittent data were obtained. The F4 satellite was located with ascending node at 2210 hours and descending node at 1010 hours geographic local time. Due to the offset between the magnetic and geographic axes, the two satellites sampled a region of magnetic local time approximately one hour to either side of the geographic local times given above while they were at low latitudes and a slightly wider range of magnetic local times at higher latitudes.

3. DATA

The total ion or electron density data were compiled for the period 1 January 1979 to 31 December 1979. This represents 12 months of F2 data and 6 months of F4 data. Data were available for 95 percent to 98 percent of each day of data collection. Data were collected in a format usable for this study on 90 percent of the possible days. The data were first compiled as 5-sec averages of the seven samples per second values of \log_{10} (density [cm⁻³]). Because of uncertainty in knowing the density when the current to the sensor reached saturation, the few data points equal to this maximum value have been discarded from the database. Because the ion density measurements are not affected by the spacecraft potential, almost all of the data come from the ion sensor. In addition, the root-mean-square (rms) of \log_{10} (density) for each 5-sec period was also calculated and saved. These two values, plus a time tag and ephemeris information, were saved for further processing.

Since the data are given and used in terms of $\log_{10}(\text{density})$ instead of density, the average values represent the median of the density (a geometric mean). A median value is generally best in comparing future data sets to past results when the underlying process is normally quite variable. For regions where the distribution of values covers a small range, the geometric mean is almost equivalent to the arithmetic mean. In cases where there is a wide variation in density (up to two orders of magnitude in the polar regions), the geometric mean is less than the arithmetic mean.

Because the ionosphere is strongly controlled by the geomagnetic field, a coordinate system based on the magnetic field should be used to represent the data. We have chosen to use the corrected geomagnetic coordinate system of Gustafsson 10, 11 for epoch 1965. We have done this by taking the geographic latitude and longitude of the satellite for zero altitude and using Gustaffson's table. This avoids the discontinuity in the invariant latitude system as the satellite crosses the equator and also avoid problems with the dip-latitude system near the poles.

The next step consisted of gathering the data into a series of statistical bins. Into each bin is placed the average value of \log_{10} (density), or median density, if the rms of the 5-sec average was less than a given value. The data were first sorted into bins representing: (a) 3° of magnetic latitude, (b) 0.25 hour of magnetic time, (c) 60° of geographic longitude, (d) 3 months of the year, and (e) each satellite. From this survey, we concluded that the variation in magnetic local time was not significant enough to justify the small divions, but the seasonal variation was not adequately described. Thus we re-generated our survey into bins representing: (a) 60° of geographic longitude, (b) 3° of corrected geomagnetic latitude, (c) 9 or 10 days, (d) either AM or PM magnetic local time and (e) each satellite for a total of 6 by 60 by 40 bins by 2 bins by 2 bins.

At first, we used a criterion of rms of the 5-sec average of $\log_{10}(\text{density}) \leq 0.02$, for all locations when selecting data to be included in a bin. This procedure eliminates occasional bad data points and data in regions with sharp density gradients and/or disturbed plasma. In particular, we wished to eliminate regions of equatorial bubbles, which are associated with spread-F, where the density can decrease by a few percent to a factor of 10^2 over a short distance and significant, small-scale density variations in the high-latitude regions. By keeping track of the percentage

Gustafsson, G. (1970) A revised corrected geomagnetic coordinate system, <u>Arkiv fur Geophysik</u>, 5:595-616.

Gustafsson, G. (1984) Corrected geomagnetic coordinates for epoch 1980, Magnetospheric Currents, T.A. Potemra, Ed., Am. Geophys. U., Washington, D.C.

of data discarded by this criteria, this survey was used to look for the longitude, latitude and seasonal variations of the occurrence of the low latitude and midlatitude density irregularities.

While the above criterion for selecting 5-sec averages was good for surveying the occurrence of low and mid-latitude density irregularities, it was not good for a global survey of the median plasma density. Much of the high latitude 5-sec averages were eliminated from the survey when the criterion of 0.02 for deviations was used. Specifically, the probability of occurrence of irregular plasma during a 5-sec sampling period using this criterion at high latitudes was always in the range of 40 percent to 100 percent. By comparison, Clark and Raitt¹² found a probability of 20 percent to 30 percent for 0.2-sec sampling periods using a criterion of the rms of the density exceeding 10 percent (which is approximately equal to a rms of \log_{10} (density) of 0.04). By using a criterion of the rms of the density exceeding 5 percent Clark and Raitt¹³ found probabilities of occurrence of 20 percent to 60 percent in the high latitude regions.

For a survey of the median plasma density, we did a separate survey using the criteria of:

rms of 5-sec sample of
$$\log_{10}(\text{density}) \le 0.02 \text{ for } -30^{\circ} \text{ to } +30^{\circ} \le 0.10 \text{ for } -60^{\circ} \text{ to } -30^{\circ}$$
 and $+30^{\circ} \text{ to } +60^{\circ} = 0.50 \text{ for } -90^{\circ} \text{ to } -60^{\circ} = 0.50 \text{ for } -90^{\circ} \text{ to } +90^{\circ} = 0.50 \text{ for } -90^{\circ} = 0.50 \text{ f$

The low latitude criterion is the same sample as that used in the irregularity survey and the high latitude criterion is chosen to be so large that almost all data not effected by telemetry errors are selected. The mid-latitude criterion was chosen to allow a transition between the high and low latitude regimes.

Clark, D. H., and Raitt, W. J. (1975) Characteristics of the high-latitude ionospheric irregularity boundary, as monitored by the total ion current probe on ESRO-4, <u>Planet. Space Sci.</u>, <u>23</u>:1643-1647.

^{13.} Clark, D. H., and Raitt, W. J. (1976) The global morphology of irregularities in the topside ionospheric as measured by the total ion current probe on ESRO-4, Planet. Space Sci., 24:873-881.

4. MEDIAN ION DENSITY

Ion density in the ionosphere as a function of latitude, longitude, time of day, season, and solar activity is given by models such as ITS-78 (Barghausen 14) and IRI-79 (Rawer 2). The ITS-78 model was based almost exclusively on measurements made from the ground up to the maximum density of the F-region at 350 km to 600 km altitude. The IRI-79 model has used topside density data from the Alouette topside sounder. However, these data were weighted more heavily for measurements near f_0 F2 than for high altitudes. As a result, it is not obvious the model densities at 830 km altitude should match the observed densities.

Our first task was to compare the measured median density with the density from previous models to determine if the previous models were accurate enough for future work in the topside ionosphere. In Figures 1 through 4 we have plotted the difference of $\log_{10}(\text{density})$ between the DMSP median density and the ITS-78 model density. Only one longitude sector is shown as an example. For some areas the difference is less than 0.3 (a factor of two in density) which is acceptable for much work, but the difference is 1.0 (a factor of ten in density) or more in other areas. In particular, the ITS-78 model is very poor in the mid-latitude region of the winter hemisphere. We have not made as thorough of a comparison of the observed densities with the IRI-79 model, but we have similar differences. As a result, we have proceeded to create a complete survey of the region of space observed by the DMSP satellites for 1979.

Figures 1 through 4 and similar representations of the data set showed that the small bin size for magnetic local time were not a necessary element of the analysis. For the morning or evening sector, the variations in \log_{10} (density) are generally less than 0.3 (a factor of two in density). In addition, there are large sections of the magnetic latitude vs magnetic local time grid in Figures 1 through 4 that have no data (the unshaded area). Thus, our further analysis is for only two local time zones (local times < 12 hours, denoted as AM, and local times > 12 hours, denoted as PM).

Bargha isen, A. F., Finney, J. W., Proctor, L. L., and Schultz, L. D. (1969)
 Predicting long-term operational parameters of high-frequency sky-wave
 telecommunication systems, ESSA Technical Rep. ERL 110-ITS 78,
 U.S. Govt. Printing Office, Washington, D.C.

LOG(N_{DMSP}) - LOG(N_{ITS78}) SUMMER 1979 GEOGRAPHIC LONGITUDE = 0° TO 60° E

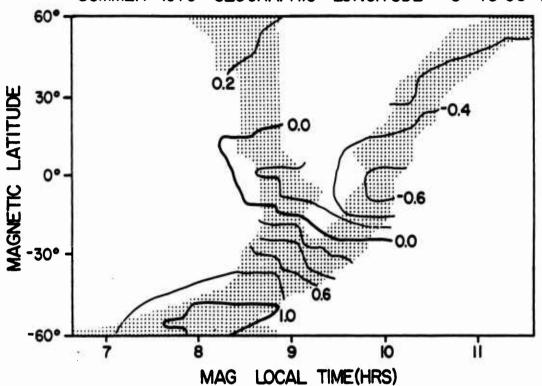


Figure 1. Difference Between Average of Log J (density), or Median of the the Density, as Given by the DMSP Data Set and the Model Value Given by ITS-78 for the Northern Hemisphere "Summer" Season (7 May to 6 August) as a Function of Magnetic Local Time in the Morning Hours and Corrected Geomagnetic Latitude for Geographic Longitudes Between 0° and 60°E

The values of \log_{10} (density) were collected, as described in Section 3, for the global survey of the median plasma density. The three levels of the rms of \log_{10} (density) which were used were chosen to eliminate the density decreases related to equatorial bubbles, the sharp gradients associated with the walls of extremely deep mid-latitude troughs, the sharp density gradients associated with local, intense auroral precipitation zones and with occasional bad points due to errors in the telemetry.

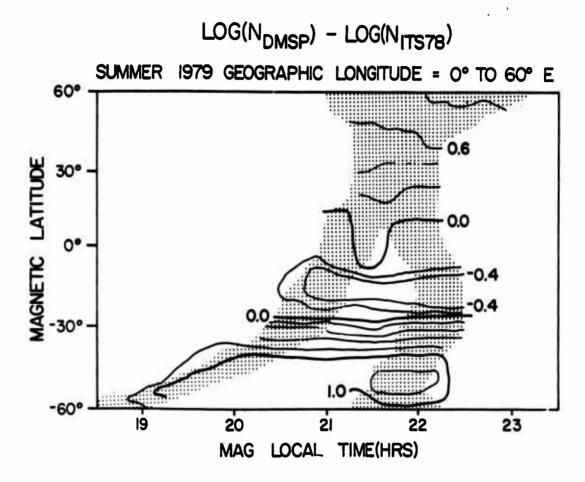


Figure 2. Same as Figure 1 Except for Magnetic Local Time in the Evening Hours

The values of \log_{10} (density) were analyzed as a function of longitude and found to have no systematic variation with longitude within a factor \pm 0.3 in \log_{10} (density) or a factor of two in density. Thus, we have grouped our data for all longitudes together into a single bin for the creation of our survey of the median density. The result is a set of two matrices of average \log_{10} (density), or median density, vs magnetic latitude and season. Figures 5 and 6 represent the matrices for the AM and PM time zones, respectively. As expected, the density near the equator tends to be greater than the density at mid- and high latitudes. The maximum values represented in Figure 6 near the equator in March may be slightly less than the true value because the ionospheric density occasionally exceeded the limits of the instrument and these points were removed from the data before the median was calculated.

Figure 3. Same as Figure 1 Except for the Northern Hemisphere "Winter" Season (7 November to 6 February)

MAG

LOCAL TIME(HRS)

At mid- and high latitudes, the density is lower in the winter season (November to February for the Northern Hemisphere and May to August in the Southern Hemisphere) than in the summer season. The winter density in the Southern Hemisphere (May to August), low density region in Figures 5 and 6 is lower than the winter density in the Northern Hemisphere (November to February). This is the same result as reported by Sojka et al 15 except that the region of low density on average

Sojka, J.J., Raitt, W.J., Schunk, R.W., Parish, J.L., and Rich, F.J. (1985) Diurnal variations of the dayside, ionospheric, mid-latitude trough in the Southern Hemisphere at 800 km; model and measurement comparison, Planet. Space Sci., 33:(No. 12):1375-1382.

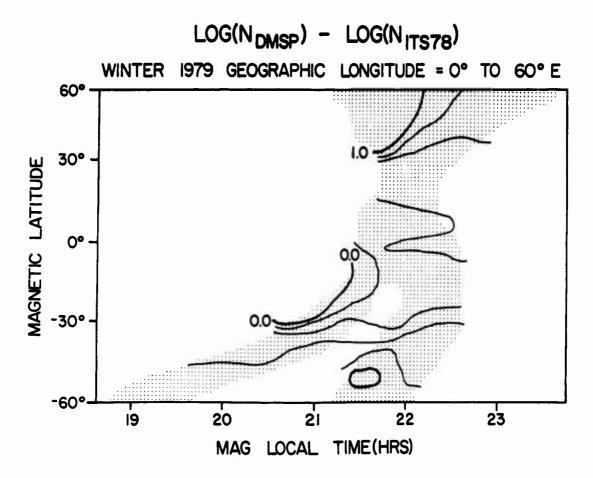


Figure 4. Same as Figure 3 Except for Magnetic Local Time in the Evening Hours

extends to lower latitude than previously noted. This implies a connection between the high and mid-latitude ionosphere regions since the lower density reported by Sojka et al¹⁵ is due to a combination of the high latitude convection pattern and the offset of the geomagnetic pole from the geographic pole.

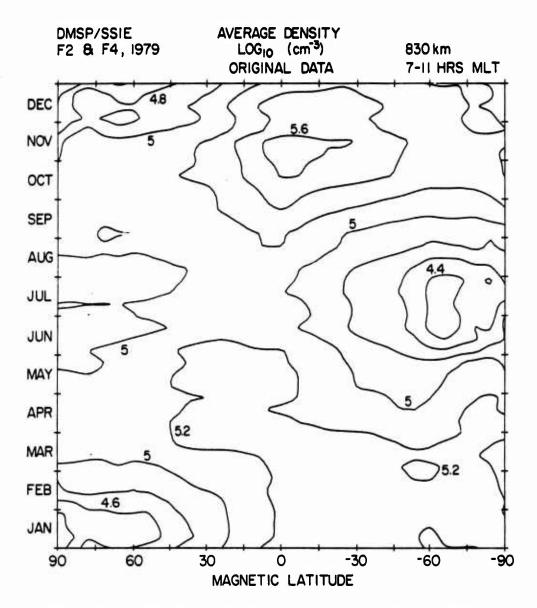


Figure 5. Contour Plot of the Average of Log₁₀ (density), or Median of the Density, During 1979, Based on Data From the SSIE Plasma Instrument on the DMSP F2 and F4 Satellites, vs Season and the Corrected Geomagnetic Latitude of the Sub-satellite Position. Data shown are for magnetic local times before 12 hours; almost all observations were made between 7 and 11 hours mag. local time

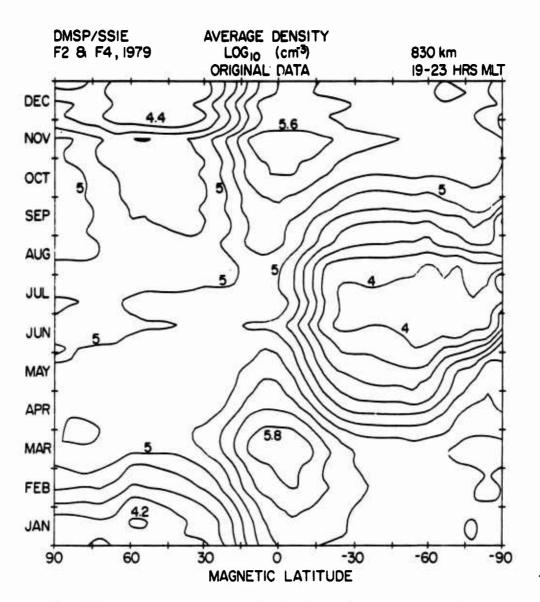


Figure 6. Same as Figure 5 Except for Data Obtained When the F2 and F4 Satellites Were in the Magnetic Local Time Regions Later Than 12 Hours

For purposes of reporting and using the results of this statistical survey, the large matrices of numbers used to generate Figure 5 and 6 are unwieldy. Thus we have fitted the values to a polynomial and we report the coefficients of the polynomial. The polynomial we have chosen is an associated Legendre polynomial of order 7 which can be expressed as:

$$\log_{10}(\text{ion density}) = \sum_{n=0}^{7} \sum_{m=-n}^{n} C_n^m b_n^m Y_n^m$$
 (2)

where

x = corrected geomagnetic latitude/90°,

y = Day Number * $2\pi/365$,

 $Y_n^m = \cos(m y) P_n^m$ for m = (0, n),

= $\sin (m y) P_n^m$ for m = (-n, -1),

C_n^m = fitting coefficient,

b_n = normalization factor, and

P_n^m = associated Legendre polynomial.

The functions b_n^m and P_n^m are defined as:

$$b_{n}^{0} = \frac{(2n)!}{2^{n}(n!)^{2}} [(2n+1)/4\pi]^{1/2}$$

$$b_{n}^{m} = \frac{(2n)!n!}{2^{n}(n!)^{2}(n-m)} [\frac{(2n+1)(n-m)!}{2\pi(n+m)!}]^{1/2}$$

$$P_{0}^{0} = 1$$

$$P_{n}^{n} = (1-x^{2})^{1/2} P_{n-1}^{n-1}$$

$$P_{n}^{m} = x P_{n-1}^{m} - R_{n}^{m} P_{n-2}^{m}$$
(3)

where

$$R_n^m = \frac{(n-1)^2 - m^2}{(2n-1)(2n-3)}, R_n^{n-1} = 0.$$

The fitting coefficients are given from the parameters COEF(N, M) in Table 1 by the relationship:

$$COEF(N, M) = C_{N-1}^{M-1} \quad \text{for } N \ge M$$
 (4)

$$COEF(N, M) = C_{M-1}^{-N}$$
 for $N < M$.

Table 1. Coefficients for Associated Legendre Polynomial Fit for Log $_{10}$ (Ion Density) at 830 km

	80		-0. 002454	9. 029328	0.003327	-0.001790	-0. 007405	-0.0:0536	-0. 006067	-0.038642		∞		-0.02850	0.01280	0.03202	0.00147	-0.01218	0. 02077	0. 00953	-0. 03013
			<u>.</u>	Ċ.	Ö	6-	-0	-0	-0.	o-				-0	0	0	0	-0	0	0	-0
	2		-0.04666	-0.01387	-0.04484	-0.04023	-0.01735	-0, 00600	0.00855	0.00306		2		-0.33832	-0.04391	-0.03728	-0.02550	-0.00795	0.01566	0.05602	0.00526
	9		-0.03680	-0.01148	0.00615	-0.00497	-0.02389	0.01051	0.00340	0. 00203		9		0.04058	0.06232	0.04575	-0.00460	0.00837	0.04867	0.00610	0.01160
	S		0.20473	-0.08758	-0.04527	-0.06148	-0.01129	0.01876	-0.00052	0.01138		ល		0.48093	-0.10546	-0.04196	-0.05730	-0.01396	0.01330	0.00419	+0.01929
AM Coefficients COEF(N, M)	4		-0.03417	0.07744	-0.01747	0.00855	-0.03367	0.01520	-0.01780	0.00207	PM Coefficients COEF(N, M)	4		-0.03864	0.04528	0.00102	-0.06514	-0.03111	-0.01608	-0.00819	0.00512
AN	က		-0.30590	-0.05355	-0, 21682	-0.06235	-0.02731	-0.00918	-0.05050	0,00260	PM	က		-0.31725	0.02902	-0.04589	-0.17119	-0.03227	-0.02161	-0.02955	0.01200
	8		-0.11666	-0.05314	-0.23725	-0.00038	-0.18805	0.03136	-0.08485	0.00100		8		-0.01955	0.24486	-0.47167	0.07978	-0.19248	0.02942	-0.19021	0.02623
	1		18, 03569	0.31439	- 0.51140	0.04421	0.01048	0.02719	- 0.05646	- 0.00866		1		17.54728	0.58212	- 1.07496	0.00776	0.16790	0.08711	- 0.14796	- 0.01168
	z											z									
		Σ	-	8	က	4	2	9	7	∞			×	-	2	က	4	သ	9	7	∞

Figures 7 and 8 represent the results of the fitting polynomial in the same format as Figures 5 and 6 for the original binned medians. Clearly all of the major components of the distribution shown in Figures 5 and 6 are reproduced in Figures 7 and 8.

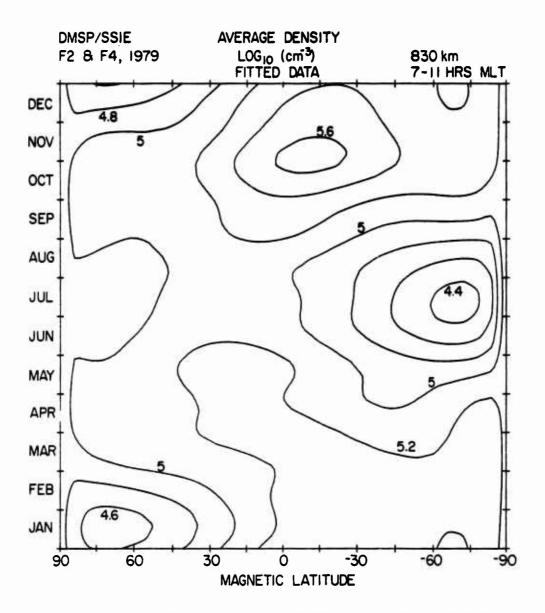


Figure 7. Contour Plot of the Polynomial Fit to the AM Density Data Shown in Figure 5

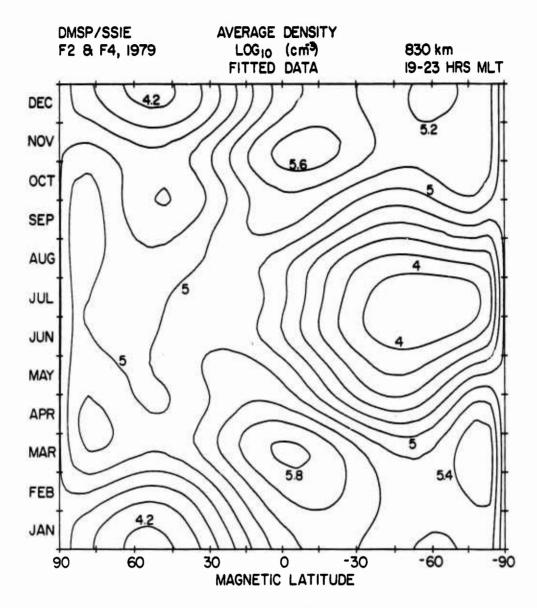


Figure 8. Contour Plot of the Polynomial Fit to the PM Density Data Shown in Figure 6

The same fitted distributions of density as shown in Figures 7 and 8 are shown again in Figures 9 and 10 in polar coordinates. Figures 9a and 10a are for the Northern Hemisphere and Figures 9b and 10b are for the Southern Hemisphere. We have re-plotted the density model to make the variation with latitude and season easier to see. One will also note that there is some failure to fit the original medians near the poles due to the coordinate system we have chosen. For the fitting procedure, the coordinates of season = January and magnetic latitude = 90° is

identical to the coordinate of July and 90°, which is logically incorrect. The result is that the fitted values are invalid between 80° (or -80°) and the pole. Between 60° and 80° (or -60° and -80°), the fitted value is not as accurate as the original data and should be used only with great care.

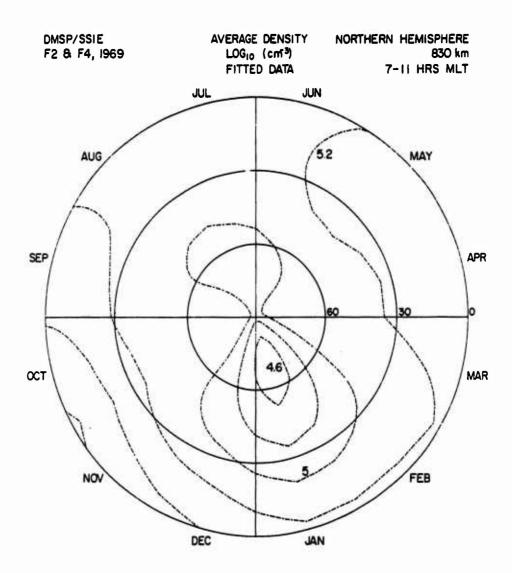


Figure 9a. Same AM Density Contour Plot as Shown in Figure 7 Except Polar Coordinates for the Northern Hemisphere

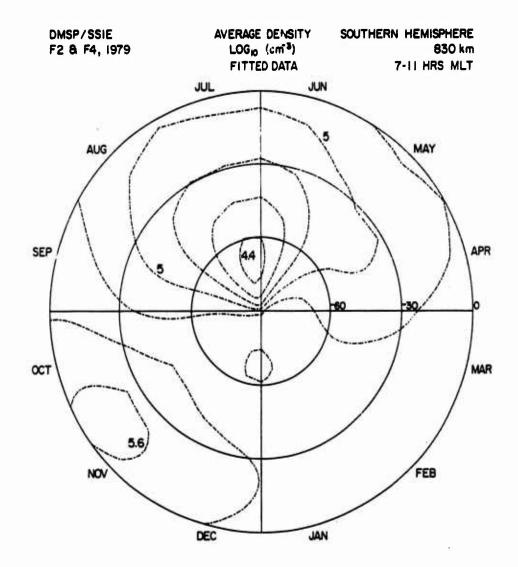


Figure 9b. Same AM Density Contour Plot as Shown in Figure 7 Except Polar Coordinates for the Southern Hemisphere

Only on a few occasions has the DMSP spacecraft potential allowed mass data to be obtained with the RPA. For F2, this has occurred only when the satellite is in darkness. For F4, we were able to make some special studies where the potential of the instrument could be biased away from the satellite potential and toward the plasma potential. On these occasions, the density at all latitudes in the AM sector is composed for > 90 percent 0+ ions. In the PM sector during the winter season, we find that light ions (H+ and He+) make up 10 percent to 60 percent of the total density in the mid-latitude regions of the low density. While we

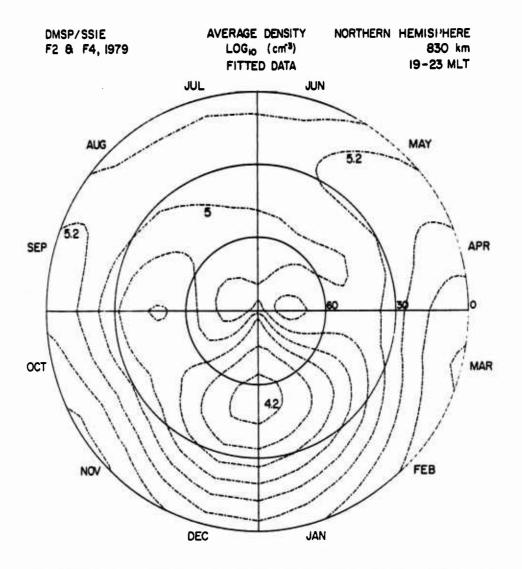


Figure 10a. Same AM Density Contour Plot as Shown in Figure 8 Except in Polar Coordinates for the Northern Hemisphere

cannot conclusively distinguish the H+ from the He+ ions, our data indicate that the majority of the light ions observed in the winter are He+. This mass distribution is consistent with the "helium bulge" which has been previously reported. A systematic survey of the available mass distribution data from the DMSP plasma instrument has not been made. Measurements from other satellites and ionospheric models generally indicate that the ionosphere at 830 km altitude is dominated by 0+ ions.

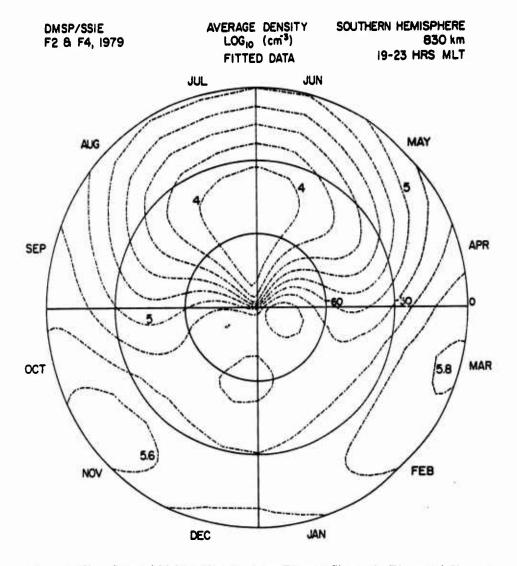


Figure 10b. Same AM Density Contour Plot as Shown in Figure 8 Except in Polar Coordinates for the Southern Hemisphere

5. LOW AND MID-LATITUDE PLASMA DENSITY IRREGULARITIES

An irregular plasma is any plasma which has small scale density gradients. The definition of small scale depends on the system being studied. At low and mid latitudes, the ionospheric plasma generally has scale lengths of the density gradients of the order of 100 to 1,000 km except in regions of turbulent plasma. By selecting plasma measurements by the criteria of the rms of a 5-sec (37 km) average of \log_{10} (density), or median of the density, being greater than 0.02, we

can find regions of plasma gradients smaller than 100 km. Then by arranging the selected "irregular" plasma measurements according to season, geographic longitude and magnetic latitude, we can find the probability of encountering the "irregular" plasma at a given time and location. The "irregular" plasma can in turn be related to degradation of radio propagation, especially in the equatorial region (Basu et al¹⁶ and DasGupta et al¹⁷).

As might be expected, the probability of encountering "irregular" plasma density regions at low and mid-latitudes in the morning sector is low (less than 10 percent of the observation) for all longitudes and all seasons. In the evening sector, there is a significant probability at certain locations and seasons. Figures 11 and 12 show the probability of encountering plasma that is "irregular" according to the rms of the 5-sec average of \log_{10} (density) > 0.02 criteria for F2 and F4 respectively during the Fall season of 1979 (7 August to 6 November). There are two zones of "irregular" plasma clearly evident: (a) near $\pm 60^{\circ}$ magnetic latitude, and (b) within 20° of the magnetic equator. The first region is mostly an extension to lower latitudes of the high latitude plasma region were the probability of rms exceeding the value of 0.02 is always very high. In addition, plasma irregularities are common on the equatorward edge (magnetic latitude = 45° to 60°) of the evening, mid-latitude trough (Basu¹⁸).

The second region is probably related to plasma bubbles which cause the "spread-F" signatures in ionograms (Hanson et al¹⁹). These bubbles are regions of low density plasma that are created near the steep bottomside gradient of the F-region shortly after sunset due to the Rayleigh-Taylor instability. The bubbles then drift upward to the altitude of the DMSP satellite and beyond. Typical dimensions for these bubbles are approximately 100 to 200 km in the east-west and vertical directions and 500 km to a few thousand km in the north-south direction along magnetic field lines. In fact, as the bubbles drift upward, they elongate in the north-south dimension.

^{16.} Basu, S., Khan, B.K., and Basu, S. (1976) Model of equatorial scintillation from in situ measurements, Radio Sci., 11:821:832.

^{17.} DasGupta, A., Aarons, J., Klobuchar, J.A., Basu, S., and Bushby, A. (1982)
Ionospheric electron content depletions associated with amplitude scintillations in the equatorial region, Geophys. Res., Letters, 9(No. 2):147-150.

Basu, S. (1978) Ogo 6 observations of small-scale irregularity structures associated with subtrough density gradients, <u>J. Geophys. Res.</u>, 83(No. A1):182-190.

^{19.} Hanson, W. B., McClure, J. P., and Sterling, D. L. (1973) On the cause of equatorial spread F, J. Geophys. Res., 78:2353-2356.

DMSP/F2 SSIE 1979 FALL PERCENTAGE OF OCCURENCE OF IRREGULARITY IN EVENING SECTOR

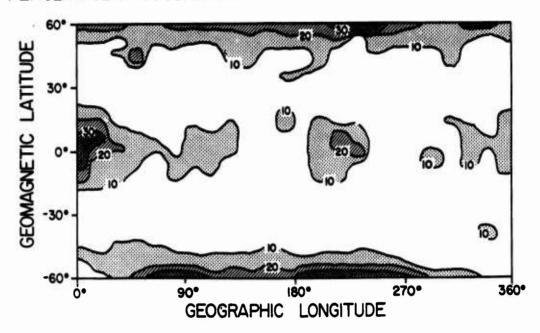


Figure 11. Contour Plot of the Percentage of F2 Data Collected in the Evening Sector (magnetic local time > 12 hours) During the Fall Season of 1979 for Which the rms of Log₁₀ (density) for a 5-sec Period Exceeded 0.02. This criteria indicates the locations of disturbed plasma

In many cases, the identification of plasma bubbles can be confirmed because the plasma density inside one of these bubbles may be less than the density in the surrounding region by as little as 20 percent or as large as a factor of 10³. This yields a significant plasma density gradient on the edges of the bubble. In addition, the plasma in and near the bubble tends to be very turbulent; this turbulence also yields significant plasma density gradients. An investigation of the longitudinal and local time distribution of obvious plasma bubbles has been presented by Young et al. In many other cases, the identification of plasma bubbles cannot be confirmed by a drop in the 5-sec average of \log_{10} (density). However, since these disturbances have the same latitudinal and local time distributions, it is reasonable to assume that these disturbed plasma regions are either the tops or sides of plasma bubbles. In proof of this assumption, we can refer to one case when the F2 and F4 satellites were near Ascension Island nearly simultaneously with the passage of the Atmospheric Explorer E satellite and when simultaneous ground observations were being made. The F2, AE-E and ground observations clearly

DMSP/F4 SSIE 1979 FALL PERCENTAGE OF OCCURENCE OF IRREGULARITY IN EVENING SECTOR

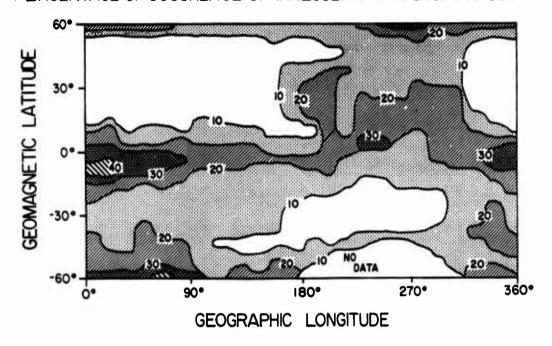


Figure 12. Same as Figure 10 Except for F4 Data Set During the Fall Season of 1979

confirmed the plasma bubbles. The F4 observations indicated large regions of disturbed plasma, but few observations of clear density depletions. The explanations for the F4 observations is that F4 was travelling nearly parallel to the magnetic meridian and passed along the side of a bubble. F2 and AE-E travelled across magnetic meridians and passed through the bubbles. As a result, we assume that most plasma disturbances observed within 20° of the magnetic equator are related to the spread-F plasma regions.

In comparing Figures 11 and 12, there is clearly a greater probability for the F4 satellite to have encountered the "irregular" plasma regions near the equator than for the F2 satellite. The explanation is that the F2 satellite is closer to the terminator and that less of the "spread-F" plasma turbulence has had enough time to be created and to drift up to the altitude of the satellite. In fact, Young et al found that as F2 drifted slowly to later local times between 1977 and 1979, the probability of encountering regions of irregular plasma progressively increased.

Previous investigations have indicated a seasonal control of the occurrence of low latitude plasma irregularities. For example, Aaron et al 20 used radio scintillation data and found a minimum in the probability of occurrence during the summer (May to July) and winter (November to January) seasons. This survey of plasma irregularities has also found that there is a seasonal dependence to the probability of encountering low latitude regions of disturbed plasma. The probability is greatest in the Spring and Fall (Figures 11 and 12), but there is a significant probability of encountering regions of irregular plasma in all seasons. Figures 13 and 14 represent a summary of the probabilities for the entire data set. For F4 the probability is weighted toward the Fall season because of a lack of data in the first six months of the year.

DMSP/F2 SSIE 1979 SUMMARY PERCENTAGE OF OCCURENCE OF IRREGULARITY IN EVENING SECTOR

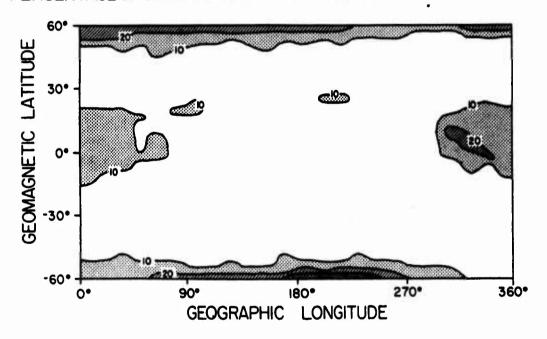


Figure 13. Same as Figure 10 Except for F2 Data Set During All Seasons of 1979

Aarons, J., Mullen, J.P., Koster, J.P., DaSilva, R.F., Medeiros, J.R., Medeiros, R.T., Bushby, A., Pantoja, J., Lanat, J., and Paulson, M.R. (1980) Seasonal and geomagnetic control of equatorial scintillations in two longitude sectors, J. Atmos. Terr. Phys., 42:861-866.

DMSP/F4 SSIE 1979 SUMMARY PERCENTAGE OF OCCURENCE OF IRREGULARITY IN EVENING SECTOR

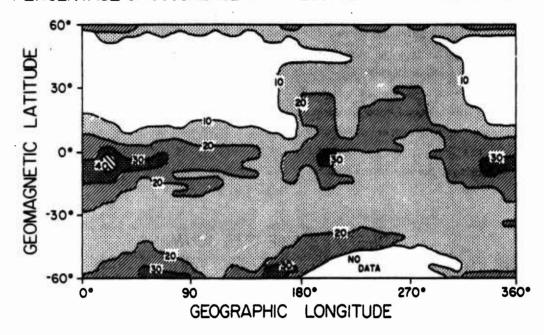


Figure 14. Same as Figure 10 Except for F4 Data Set During All Seasons of 1979

One of the most striking features of the low latitude, plasma irregularity distribution is the longitudinal dependence. The probability is always highest in a zone between 340° and 60°E longitude and a second zone between 195° and 240°E longitude. Young et al⁹ offered the suggestion for F2 that this longitude dependence was related to the angle between the satellite motion and the magnetic meridian. If this were true, then the longitudinal dependence for F4 would be significantly different, which it is not. This longitudinal dependence is very similar to that found by Basu et al 16 using OGO-6 data during November and December of 1969 and 1970 and by Maruyama and Matuura using ISS-b data from 1978 to 1979. This includes a secondary maximum in the probability of irregular plasma near 180° longitude which Young et al⁹ failed to find in their limited set of DMSP data. The ISS-b results also indicate the presence of a secondary maximum in the irregular plasma near 180° longitude, but only for August and September 1979, which is consistent with our results. The criterion used by Basu et al 16 for locating a region of plasma irregularity was $\Delta N > 10^4$ cm⁻³ in any part of a 3° latitude bin which is a more liberal definition than that used herein or used by Clark and Raitt. 12, 13 The use of ΔN instead of $\Delta N/N$ is that only regions with plasma

irregularities strong enough to locally create scintillation of radio signals are counted. The DMSP observations are made at an altitude above the plasma that disrupts the radio signal. However, from a practical point of view, the difference is of minor importance at low latitudes. If a disturbed plasma is observed over the equatorial region by DMSP, then the plasma in the column containing the satellite is disturbed enough to disrupt radio signals. The difference in percentages between the results of the Basu study and the one herein also represents the fact that not all plasma disturbances observed at 400 km in the 19 to 23 hour sector have had enough time to drift up to 830 km altitude.

6. DISCUSSION AND CONCLUSIONS

The density survey given herein is based on the largest data set of any model for the topside ionosphere. Since standard models of the ionosphere do not accurately represent the topside ionosphere, the survey given herein plus similar data sets from other programs should be included in future revisions of the standard models. Changes in the topside of the standard ionospheric models such as indicated herein will have only a small effect (approximately 25 percent or less) on the total electron content (TEC) in a column through the ionosphere (cm⁻²), which can be measured directly with measurements on the ground. However, McNamara and Wilkinson²¹ have indicated that just such a change in the IRI model is needed to improve the agreement of the TEC values from the model with TEC measurements at 30°S. In addition, there are applications where TEC is estimated from topside measurements combined with profiles from the standard models of the ionosphere. For these applications, it is critical to have the standard models fit the measurements given herein.

The limitations of this survey are that it only covers one altitude (830 km), two local times (approximately 1000 hours and 2200 hours) and one point on the solar cycle (solar maximum). We looked for evidence of a geomagnetic activity effect, using the Kp index, on the density and did not find a statistically meaningful one. The altitude dependence of the density can be partially determined from other satellites, such as ISS-b, and partially from temperature measurements made by plasma sensors on future DMSP satellites. Some additional local times can be surveyed by future DMSP satellites, but the time sectors 1200 to 1700 and 0000 to 0500 will have to be filled in with data from other programs since DMSP satellites are never placed in those sectors.

McNamara, L. F., and Wilkinson, P. J. (1983) Predictions of total electron content using International Reference Ionosphere, J. Atmos. Terr. Phys., 45(No. 2/3):169-174.

The solar cycle is known to have a significant effect on the bottomside of the ionosphere and on the temperature of the ionosphere. Thus, it is reasonable to assume that the median densities given herein would not be accurate near the minimum of the solar cycle. To illustrate this solar cycle variation, we present in Table 2 a comparison of some median densities observed by DMSP in 1979 and by a very similar plasma instrument on the HILAT satellite in 1984, which is close to the time of solar minimum. The HILAT satellite is also a polar orbiting satellite near 830 km altitude, but only high latitude data are available from it. At 60° magnetic latitude, where the daily variations of plasma density are less than at higher latitudes, the median density changes by a factor of four to ten as a result of the different phases of the solar cycle. Future DMSP data sets can provide more information about the dependence on the solar cycle of the median plasma density in the topside ionosphere.

The results of the density survey given herein compare well with the results of ESRO 4 (Clark et al²²) obtained at 600 to 700 km altitude during October 1968 to April 1969, which was also near the maximum in the solar cycle. The implication is that the density is approximately the same at the same phase of two different solar cycles. The report of the ESRO 4 data does not show any seasonal dependence, but does show a local time dependence over 24 hours. Their results indicate that there are not strong gradients of density with respect to local time for the local time sectors observed by DMSP. However, the results from ESRO 4 did show a strong local time gradient in the 1400 to 1800 hour and 0000 to 0400 hour time sectors.

The median plasma densities given herein also compare well with the results of ISS-b (Hakura³) obtained at 1100 km altitude during 22 October 1978 to 22 February 1979 and during 22 April 1979 to 22 August 1979 at 1000 and 2200 hours local time. Unlike this survey, the results from ISS-b include the densities of H+, He+ and O+ separately. These results clearly show that He+ is often the dominant ion at mid-latitudes of the winter, nighttime ionosphere. Unlike the results of this study, Hakura has found a significant longitudinal dependence to the density, especially in the winter, low-density region. Since their results cover approximately the same time period as our results, we speculate that the difference is due either to the result of temperature gradients which produce different scale heights or simply due to different sampling and collation methods.

^{22.} Clark, D. H., Raitt, W. J., and Willmore, A. P. (1972) The global morphology of electron temperature in the topside ionosphere as measured by an a. c. Langmuir probe, J. Atmos. Terr. Phys., 34:1865-1888.

Table 2. Comparison of Median Density at 60° Magnetic Latitude Observed at 830 km Altitude by DMSP in 1979 and HILAT in 1984

	(Density - 10 ⁴ cm ⁻³)							
	Sum	mer	Winter					
	DMSP	HILAT	DMSP	HILAT				
Morning 07 to 11 Hrs MLT	10	1. 8	2. 5	0. 18				
Evening 19 to 23 HRS MLT	10	2.5	4 to 7	0.36				

The low latitude plasma irregularity results given herein correlate well with those of Basu et al 16 which can be related to radio scintillations. This demonstrates that similar DMSP data can be used to locate regions where transionospheric HF, VHF and UHF radio signals will be degraded by the ionospheric plasma. In the equatorial regions, finding a region of disturbed plasma will be a sufficient indicator of radio degradation, but not a necessary one since as many as half of the regions of disturbed plasma have not drifted up to the altitude of the DMSP satellite.

In the high latitude regions, finding regions of disturbed plasma must be combined with estimates of the plasma density at the peak of the F-region (approximately 350 km altitude) in order to have a necessary condition for radio degradation. Low density plasma cannot disturb radio signals unless the degree of the degree of irregularity in the density ($\Delta N/N$) increases. A survey of the total electron content (TEC) made simultaneously with the HILAT measurements in Table 2 (S. Basu, priv. comm.) showed a change of only a factor of two between Winter 1984 and Summer 1984. This implies a much smaller change in the density at the peak of the F-region than was observed in the topside ionosphere at the satellite altitude. This implies a seasonal variation of the temperature in the topside ionosphere.

In order to relate high latitude, topside plasma irregularity to scintillations, more information is needed about the seasonal and solar cycle dependence of temperature in the topside. In addition, the altitude dependence of $\Delta N/N$ needs to be investigated. It is often assumed that $\Delta N/N$ is independent of altitude but there is little proof for or against that assumption.

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